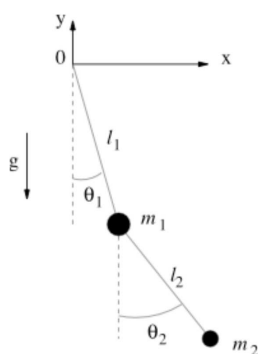


Due *in class* Thursday February 28th*Lagrangian Mechanics:*

1. Two blocks, both of mass M , are connected with a spring with spring constant k . They slide in the x -direction on a frictionless surface.
 - (a) Use the Lagrangian formalism to find coupled differential equations for the positions x_1 and x_2 of the two blocks.
 - (b) Use these equations to show that the center-of-mass of the system moves with constant velocity.
 - (c) Find the period of oscillation of the two masses with respect to each other.
2. Consider the *double pendulum* in the figure below, consisting of two bobs of mass m_1 and m_2 at the ends of massless rods of length l_1 and l_2 in which the second rod swings from the first mass.
 - (a) Derive (but do not solve) the equations of motion for the two bobs in terms of their coordinates θ_1 and θ_2 . You may find the arithmetic easier if you make use of the trigonometric identity

$$\cos A \cos B + \sin A \sin B = \cos (A - B)$$

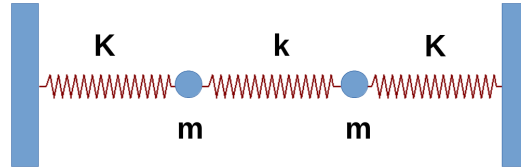
- (b) Find the simplified version of these equations for the case in which θ_1 and θ_2 are both small angles.

*Additional problems on the back of this page.*

Normal Modes of Oscillation:

3. $K\mathcal{E}K$ Problem 6.3.

4. See the figure below. Three springs and two equal masses lie between two walls. The spring constant K of the outer springs is larger than the spring constant k of the inner spring. Let x_1 and x_2 be the positions of the left and right masses, respectively, relative to their equilibrium positions.

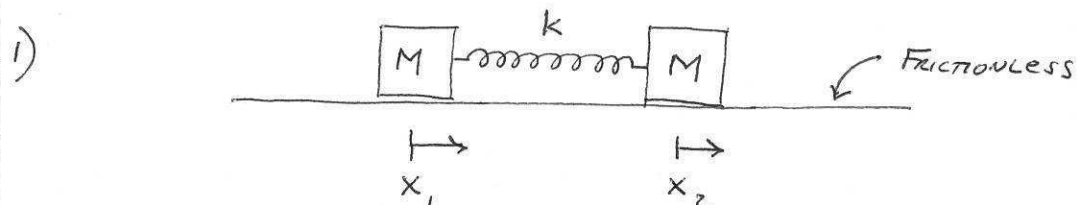


- (a) Find general oscillatory solutions for $x_1(t)$ and $x_2(t)$.
(b) Assume $K/m = 1.0/\text{s}^2$ and $k = 0.2 K$. Plot x_1 and x_2 versus time for the normal modes of the system. Assume the equilibrium distance between the two masses is 1.0 m.
(c) Now, find the particular solutions for x_1 and x_2 , for the initial conditions:

$$\begin{aligned}x_1(0) &= 0.0 \text{ m} \\ \dot{x}_1(0) &= 0.0 \text{ m/s} \\ x_2(0) &= 0.4 \text{ m} \\ \dot{x}_2(0) &= 0.0 \text{ m/s}\end{aligned}$$

Plot these solutions for t in the range $0 - 100$ s.

(Note: If you don't have access to a graphing program on your computer, have a look at online plotting tools such as **WolframAlpha** and **fooplot.com**.)



a) $L = K - U = \frac{1}{2} M \dot{x}_1^2 + \frac{1}{2} M \dot{x}_2^2 - \frac{1}{2} k (x_2 - x_1)^2$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_1} \right) - \frac{\partial L}{\partial x_1} = M \ddot{x}_1 - k(x_2 - x_1) = 0$$

$$\textcircled{1} \quad \ddot{x}_1 - \omega^2 x_2 + \omega^2 x_1 = 0$$

$$\omega = \sqrt{\frac{k}{m}}$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_2} \right) - \frac{\partial L}{\partial x_2} = M \ddot{x}_2 + k(x_2 - x_1) = 0$$

$$\textcircled{2} \quad \ddot{x}_2 + \omega^2 x_2 - \omega^2 x_1 = 0$$

b) Add Eqs $\textcircled{1}$ and $\textcircled{2}$:

$$\ddot{x}_1 + \ddot{x}_2 = 0$$

ACCELERATION

OF C.O.M. = $\frac{M \ddot{x}_1 + M \ddot{x}_2}{2M} = \frac{1}{2} (\ddot{x}_1 + \ddot{x}_2) = 0$

\therefore COM moves w/ constant \vec{v}

c) Subtract Eq $\textcircled{1} - \textcircled{2}$

$$(\ddot{x}_1 - \ddot{x}_2) + 2\omega^2 (x_1 - x_2) = 0$$

Simple harmonic oscillator in $(x_1 - x_2)$ w/ freq $\omega' = \sqrt{2} \omega$
 $= \sqrt{\frac{2k}{m}}$

Period of oscillation $T = \frac{2\pi}{\omega'} = \pi \sqrt{\frac{2m}{k}}$

2) EASIEST TO START BY WRITING DOWN LAGRANGIAN IN CARTESIAN COORDS, THEN USE TRIGONOMETRY TO CONVERT:

$$a) L = K - U = \frac{1}{2} m_1 (\dot{x}_1^2 + \dot{y}_1^2) + \frac{1}{2} m_2 (\dot{x}_2^2 + \dot{y}_2^2) - m_1 g y_1 - m_2 g y_2$$

$$\begin{aligned} \text{Now, } x_1 &= l_1 \sin \theta_1, & \dot{x}_1 &= l_1 \dot{\theta}_1 \cos \theta_1, \\ y_1 &= -l_1 \cos \theta_1, & \dot{y}_1 &= l_1 \dot{\theta}_1 \sin \theta_1, \\ x_2 &= l_1 \sin \theta_1 + l_2 \sin \theta_2, & \dot{x}_2 &= l_1 \dot{\theta}_1 \cos \theta_1 + l_2 \dot{\theta}_2 \cos \theta_2, \\ y_2 &= -l_1 \cos \theta_1 - l_2 \cos \theta_2, & \dot{y}_2 &= l_1 \dot{\theta}_1 \sin \theta_1 + l_2 \dot{\theta}_2 \sin \theta_2 \end{aligned}$$

Plug THESE INTO EXPRESSION FOR LAGRANGIAN

$$\begin{aligned} L = & \frac{1}{2} m_1 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 l_2^2 \dot{\theta}_2^2 + m_2 l_1 l_2 \dot{\theta}_1 \dot{\theta}_2 (\cos \theta_1 \cos \theta_2 + \sin \theta_1 \sin \theta_2) \\ & + m_1 g l_1 \cos \theta_1 + m_2 g l_1 \cos \theta_1 + m_2 g l_2 \cos \theta_2 \end{aligned}$$

$$\text{use } \cos A \cos B + \sin A \sin B = \cos(A - B)$$

$$\star L = \frac{1}{2} m_1 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 l_2^2 \dot{\theta}_2^2 + m_2 l_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \cos(\theta_1 - \theta_2) + m_1 g l_1 \cos \theta_1 + m_2 g l_2 \cos \theta_1 + m_2 g l_2 \cos \theta_2$$

Plug L into Euler-Lagrange Equations, for generalized coordinates θ_1, θ_2 :

$$\begin{aligned} \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}_1} \right) - \frac{\partial L}{\partial \theta_1} &= (m_1 + m_2) l_1^2 \ddot{\theta}_1 + m_2 l_1 l_2 [\ddot{\theta}_2 \cos(\theta_1 - \theta_2) + \dot{\theta}_2^2 \sin(\theta_1 - \theta_2)] \\ &+ (m_1 + m_2) g l_1 \sin \theta_1 = 0 \end{aligned}$$

$$\begin{aligned} \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}_2} \right) - \frac{\partial L}{\partial \theta_2} &= m_2 l_2^2 \ddot{\theta}_2 + m_2 l_1 l_2 [\ddot{\theta}_1 \cos(\theta_1 - \theta_2) - \dot{\theta}_1^2 \sin(\theta_1 - \theta_2)] \\ &+ m_2 g l_2 \sin \theta_2 = 0 \end{aligned}$$



2a) ... continued

rewriting ...

$$\ddot{\theta}_1 + \frac{m_2}{m_1 + m_2} \frac{l_2}{l_1} (\ddot{\theta}_2 \cos(\theta_1 - \theta_2) + \dot{\theta}_2^2 \sin(\theta_1 - \theta_2)) + \frac{g}{l_1} \sin \theta_1 = 0$$

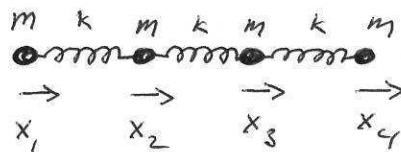
$$\ddot{\theta}_2 + \frac{l_1}{l_2} (\ddot{\theta}_1 \cos(\theta_1 - \theta_2) - \dot{\theta}_1^2 \sin(\theta_1 - \theta_2)) + \frac{g}{l_2} \sin \theta_2 = 0$$

b) Small angle approximation, $\sin \theta \approx \theta$
 $\cos \theta \approx 1$

$$\ddot{\theta}_1 + \left(\frac{m_2}{m_1 + m_2} \right) \frac{l_2}{l_1} (\ddot{\theta}_2 + \dot{\theta}_2^2 (\theta_1 - \theta_2)) + \frac{g}{l_1} \theta_1 = 0$$

$$\ddot{\theta}_2 + \frac{l_1}{l_2} (\ddot{\theta}_1 - \dot{\theta}_1^2 (\theta_1 - \theta_2)) + \frac{g}{l_2} \theta_2 = 0$$

3) K&K 6.3



$$L = K - U = \frac{1}{2} m \dot{x}_1^2 + \frac{1}{2} m \dot{x}_2^2 + \frac{1}{2} m \dot{x}_3^2 + \frac{1}{2} m \dot{x}_4^2 - \frac{1}{2} k (x_2 - x_1)^2 - \frac{1}{2} k (x_3 - x_2)^2 - \frac{1}{2} k (x_4 - x_3)^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_1} \right) - \frac{\partial L}{\partial x_1} = m \ddot{x}_1 - k (x_2 - x_1) = 0$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_2} \right) - \frac{\partial L}{\partial x_2} = m \ddot{x}_2 + k (x_2 - x_1) - k (x_3 - x_2) = 0$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_3} \right) - \frac{\partial L}{\partial x_3} = m \ddot{x}_3 + k (x_3 - x_2) - k (x_4 - x_3) = 0$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_4} \right) - \frac{\partial L}{\partial x_4} = m \ddot{x}_4 + k (x_4 - x_3) = 0$$

$$\text{Let } \omega_0 = \sqrt{\frac{k}{m}}, \quad \beta = \frac{\omega^2}{\omega_0^2}, \quad \ddot{x}_i = -\omega^2 x_i$$

THEN ABOVE CAN BE WRITTEN AS

$$\beta x_1 = (x_1 - x_2)$$

$$\beta x_2 = (2x_2 - x_1 - x_3)$$

$$\beta x_3 = (2x_3 - x_2 - x_4)$$

$$\beta x_4 = (x_4 - x_3)$$

By symmetries & condition that C.O.M. is at rest, possible normal modes are

$$\textcircled{I} \quad (x_4 = x_1) \text{ and } (x_3 = x_2)$$

$$\textcircled{II} \quad (x_4 = -x_1) \text{ and } (x_3 = -x_2)$$



CASE I Unique EQNS ARE

$$\textcircled{c} \quad \beta x_1 = x_1 - x_2$$

$$\textcircled{cc} \quad \beta x_2 = (2x_2 - x_1 - x_2) = (x_2 - x_1)$$

$$\underbrace{x_1(1-\beta)}_{\textcircled{c}} = \underbrace{x_2}_{\textcircled{cc}} = \frac{x_1}{(1-\beta)}$$

$$1 + \beta^2 - 2\beta = 1$$

$$\beta^2 - 2\beta$$

$$\beta = 2 = \frac{\omega^2}{\omega_0^2}; \quad \boxed{\omega = \sqrt{\frac{2k}{m}}}$$

CASE II Unique EQUATIONS ARE

$$\textcircled{cc} \quad \beta x_1 = x_1 - x_2$$

$$\textcircled{cu} \quad \beta x_2 = 2x_2 - x_1 + x_2 = 3x_2 - x_1$$

$$\underbrace{x_1(1-\beta)}_{\textcircled{cc}} = \underbrace{x_2}_{\textcircled{cu}} = \frac{x_1}{3-\beta}$$

$$(1-\beta)(3-\beta) = 1$$

$$\beta^2 - 4\beta + 2 = 0$$

$$\beta = 2 \pm \sqrt{2} = \frac{\omega^2}{\omega_0^2}$$

$$\therefore \boxed{\omega = \sqrt{\frac{(2 \pm \sqrt{2})k}{m}}}$$

$$4) \quad L = \frac{1}{2} m \dot{x}_1^2 + \frac{1}{2} m \dot{x}_2^2 - \frac{1}{2} K x_1^2 - \frac{1}{2} K (x_2 - x_1)^2 - \frac{1}{2} K x_2^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_1} \right) - \frac{\partial L}{\partial x_1} = m \ddot{x}_1 + K x_1 - k (x_2 - x_1) = 0$$

$$\ddot{x}_1 + \left(\frac{k+K}{m} \right) x_1 - \frac{k}{m} x_2 = 0 \quad (i)$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_2} \right) - \frac{\partial L}{\partial x_2} = m \ddot{x}_2 + K x_2 + k (x_2 - x_1) = 0$$

$$\ddot{x}_2 + \left(\frac{k+K}{m} \right) x_2 - \frac{k}{m} x_1 = 0 \quad (ii)$$

$$\text{Add } (i) + (ii) \quad (\ddot{x}_1 + \ddot{x}_2) + \frac{K}{m} (x_1 + x_2) = 0$$

$$\text{Subtract } (i) - (ii) \quad (\ddot{x}_1 - \ddot{x}_2) + \left(\frac{2k+K}{m} \right) (x_1 - x_2) = 0$$

So $x_+ = x_1 + x_2$ and $x_- = x_1 - x_2$ HAVE GENERAL SOLUTIONS

$$x_+ = A_+ \cos(\omega_+ t + \phi_+) \quad \omega_+ = \sqrt{\frac{K}{m}}$$

$$x_- = A_- \cos(\omega_- t + \phi_-) \quad \omega_- = \sqrt{\frac{2k+K}{m}}$$

$$\text{And } x_1 = \frac{x_+ + x_-}{2} = \frac{A_+}{2} \cos(\omega_+ t + \phi_+) + \frac{A_-}{2} \cos(\omega_- t + \phi_-)$$

$$x_2 = \frac{x_+ - x_-}{2} = \frac{A_+}{2} \cos(\omega_+ t + \phi_+) - \frac{A_-}{2} \cos(\omega_- t + \phi_-)$$

Apply INITIAL CONDITIONS $x_1(t=0) = 0$, $x_2(t=0) = 0.4 \text{ m}$, $\dot{x}_1(0) = \dot{x}_2(0) = 0$

\Rightarrow VELOCITY IC's give $\phi_+ = \phi_- = 0$

\Rightarrow THEN POSITION IC's $\frac{A_+}{2} + \frac{A_-}{2} = 0$

$$\frac{A_+}{2} - \frac{A_-}{2} = 0.4$$

$$\hookrightarrow \boxed{A_+ = 0.4 \text{ m}, A_- = -0.4 \text{ m}}$$

See Attached Plots \rightarrow

Normal Modes of System

